

The lower tail problem for homogeneous functionals of stable processes with no negative jumps

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Abstract. Let Z be a strictly α -stable real Lévy process ($\alpha \in (1,2]$) and X be a fluctuating β -homogeneous additive functional of Z. We investigate the asymptotics of the first passage-time of X above 1, and give a general upper bound. When Z has no negative jumps, we prove that this bound is optimal and does not depend on the homogeneity parameter β . This extends a result of Isozaki (1996) and solves partially a conjecture of Shi (2003).

1. Introduction

Let X be a real process starting from 0 and $T = \inf\{t > 0, X_t > 1\}$ be its first passage time above level one. Studying the law of T is a classical and important question in probability, with many applications especially in finance and insurance. Since it is often difficult to compute the exact distribution of T, even when X is continuous, one is sometimes interested in the asymptotic behaviour of the ruin probability:

$$\mathbb{P}\left[T > t\right], \qquad t \to \infty. \tag{1.1}$$

In a self-similar framework, this problem is equivalent to the estimate of the so-called lower tails of X:

$$\mathbb{P}\left[\sup_{t\in[0,1]}X_t<\varepsilon\right],\qquad\varepsilon\to0. \tag{1.2}$$

The latter is a less classical question, but striking connections have been made over the years between (1.2) and subjects as different as e.g. increase points for Lévy processes - see section VI.5 in Bertoin (1996), Hausdorff dimension of regular points for the inviscid Burgers equation with random initial impulse Sinai (1992b); Molchan and Khokhlov (2004), capture times of Brownian pursuits and zeroes of

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random polynomials - see Li and Shao (2004); Shao (2003) and also the references therein for other relevant issues.

Notice that despite formal resemblance, the estimate (1.2) has less to do with the so-called small deviation (or small ball) problem for X, which deals with the asymptotics

$$\mathbb{P}\left[\sup_{t\in[0,1]}|X_t|<\varepsilon\right],\qquad\varepsilon\to0.$$

Indeed the latter is usually studied, in a large deviation spirit, under the logarithmic scale, whereas in (1.2) the investigated speed of convergence speed is most of the time polynomial: one expects a behaviour of the type

$$\mathbb{P}\left[T>t\right] \; = \; t^{-\theta+o(1)}, \qquad t\to\infty, \tag{1.3}$$

for some finite positive constant θ , whose exact value is actually relevant in several problems of statistical physics - see Bhattacharya et al. (2007); Bray and Gonos (2005) and the numerous references therein. In this physical literature, the exponent θ is often mentioned as a persistence or survival exponent because in general, the probability in (1.3) behaves modulo some o(1) term like the probability that X starting from some x>0 remains positive up to time $t\to +\infty$. This expected polynomial asymptotic forces to work under the natural scale, making the problem significantly more delicate than usual small ball probability estimates and actually very few explicit exponents are known, especially in a Non-Markovian framework. So far, up to some generalisation of integrated Brownian motion which will be detailed soon afterwards, the problem of computing θ seems to have been solved only for the three following self-similar processes:

X	θ	Ref.
Integrated Brownian motion	1/4	McKean (1963); Goldman (1971)
Lévy α-stable process	ρ	Bingham (1975)
Fractional Brownian motion	1-H	Molchan (1999)

In the above, the parameter ρ stands for the positivity parameter of the Lévy α -stable process, which will be recalled henceforth, whereas the parameter H is the usual Hurst parameter of the fractional Brownian motion. Recently, two conjectures were stated on the explicit expression of θ for integrated processes, generalizing the value 1/4 for integrated Brownian motion:

X	θ	Ref.
Integrated Lévy α -stable process	$(\alpha-1)_+/2\alpha$	Shi (2003)
Integrated fractional Brownian motion	H(1-H)	Molchan and Khokhlov (2004)

In this paper we will be concerned with the first conjecture. More precisely we will consider the slightly more general situation where X is a fluctuating β -homogeneous additive functional of a strictly α -stable Lévy process Z:

$$X_t \; := \; A_t^{(eta)} \; = \; \int_0^t |Z_s|^eta {
m sgn}(Z_s) \; {
m d} s, \quad t \geq 0,$$

for $\alpha \in (1,2]$ and $\beta > -(\alpha+1)/2$. More details and the reason why the above signed integral always makes sense will be given in the next section. The process $\{A_t^{(\beta)},\ t \geq 0\}$ is $(1+\beta/\alpha)$ -self-similar, but is not Markovian and has no stationary

increments. In the case $\beta=1$, it is the area process associated with Z, in other words the integrated Lévy α -stable process:

$$A_t^{(1)} = \int_0^t Z_s \, \mathrm{d}s, \quad t \ge 0,$$

and is also an α -stable process in the general sense of Samorodnitsky and Taqqu (1994). However, in the case $\beta \neq 1$, it is no more a stable process. When $\beta = -1$, the process $A^{(\beta)}$ is up to a multiplicative constant the Hilbert transform of $\{L(t,x),x\in\mathbb{R}\}$ the local time process of Z - see section V.2 in Bertoin (1996). When $\beta<-1$, it can be viewed as a fractional derivative of $\{L(t,x),x\in\mathbb{R}\}$ - see section 2 in Fitzsimmons and Getoor (1992). Our first result is the following:

Theorem A. With the above notations, there exists a finite positive constant K such that

$$\mathbb{P}[T > t] \le \mathcal{K} t^{-(\alpha - 1)/2\alpha}, \qquad t \to +\infty. \tag{1.4}$$

Besides, the constant K is explicit when $\alpha = 2$ or $\beta = -1$.

In the Brownian case $\alpha=2$, this result had been obtained previously by Isozaki (1996) for $\beta>0$, with the help of a Wiener-Hopf factorization in two dimensions extending an identity of Spitzer-Rogozin in discrete time, and an asymptotic analysis of the Wiener-Hopf factors. We remark that these analytical arguments can be carried out to the Non-Gaussian case without much difficulty. Our main novelty consists in dealing with non necessarily symmetric stable processes.

In Isozaki (1996) it is also proved that when $\alpha=2$, the exponent $(\alpha-1)/2\alpha=1/4$ is optimal, in the sense that (1.4) holds as well with an inequality in the other direction and a positive constant $\mathcal{K}' < \mathcal{K}$. With the help of excursion theory, this result in the case $\alpha=2$ has been then significantly refined in Isozaki and Kotani (2000), where it is proved that

$$\mathbb{P}[T > t] \sim \mathcal{K}'' t^{-1/4}, \qquad t \to +\infty,$$

for some explicit constant $\mathcal{K}'' \in (\mathcal{K}', \mathcal{K})$. At a less precise level, finding the right exponent θ in (1.3) remains a tantalizing question in the Non-Gaussian case, because of the jumps of Z which make analytical arguments untractable. Results in the direction of Shi's conjecture had been given in Devulder et al. (2005), with unfortunately a serious gap in the proof. In this paper, we aim at tackling this problem in the spectrally positive situation:

Theorem B. Suppose that Z has no negative jumps. Then with the above notations, there exists a positive constant k such that

$$\mathbb{P}[T > t] \ge k t^{-(\alpha - 1)/2\alpha} (\log t)^{-1/2\alpha}, \qquad t \to +\infty.$$

Putting the two theorems together entails that $\theta = (\alpha - 1)/2\alpha$ for the processes $A^{(\beta)}$ as soon as Z has no positive jumps and, taking $\beta = 1$, solves Shi's conjecture in this particular case. It may seem surprising that the exponent θ shows no dependence on β , since the processes $A^{(\beta)}$ look different according as $\beta < 0$ or $\beta \geq 0$. Notice however that this dependence is retrieved in the lower tail formulation: by self-similarity one has

$$\mathbb{P}\left[\sup_{t\in[0,1]}A_t^{(\beta)}<\varepsilon\right] = \varepsilon^{(\alpha-1)/2(\alpha+\beta)+o(1)}, \qquad \varepsilon\to 0,$$

as soon as Z has no positive jumps.

Our proof of Theorem B is reminiscent of an argument of Sinai (1992a) who had considered this question for $\beta=1$, in discrete time. Roughly, we show that T belongs to an excursion interval of Z whose left-end can be controlled appropriately. In Sinai (1992a) - see the end of Section 3 therein, the *simplicity* assumption on the random walk is crucial to get the control. Here, we remark that this assumption can be translated in continuous time into that of the absence of negative jumps, since then X accumulates positive value between T and the next hitting time of 0 by Z. Alternatively, this accumulating argument is similar to that of Lemma 1 in Bertoin (1998), which is key in analyzing the solution of the inviscid Burgers equation whose initial data is a Lévy process with no positive jumps. The analogy is actually not very surprising, and we will give more explanation for this at the end of the paper.

To conclude this introduction, notice that Theorems A and B readily entails the following result on the positive moments of T:

Corollary. Suppose that Z has no negative jumps. Then for every k > 0,

$$k < (\alpha - 1)/2\alpha \iff \mathbb{E}[T^k] < +\infty.$$

In the third section of this paper, after proving Theorems A and B, we will give some more results in this vein concerning the relevant random variable Z_T .

2. Preliminaries

2.1. Recalls on stable processes and their homogeneous functionals. Consider $Z = \{Z_t, t \geq 0\}$ a real strictly α -stable Lévy process with index $\alpha \in (1, 2]$, viz. a process with stationary and independent increments which is $(1/\alpha)$ -self-similar:

$$\{X_{kt}, \ t \ge 0\} \stackrel{d}{=} \{k^{1/\alpha}X_t, \ t \ge 0\},$$

for all k>0. Its Lévy-Khintchine exponent $\Psi(\lambda)=-\log \mathbb{E}[e^{i\lambda Z_1}]$ is given by

$$\Psi(\lambda) = \kappa |\lambda|^{\alpha} (1 - i\chi \operatorname{sgn}(\lambda) \tan(\pi \alpha/2)), \quad \lambda \in \mathbb{R}, \tag{2.1}$$

where $\kappa > 0$ is the scaling parameter and $\chi \in [-1, 1]$ is the skewness parameter. The positivity parameter $\rho = \mathbb{P}[Z_t > 0]$, independent of t by self-similarity, is given by

$$\rho = \frac{1}{2} + (\pi \alpha)^{-1} \arctan(\chi \tan(\pi \alpha/2))$$

(see Theorem 2.6.3 in Zolotarev (1986)). Since $\alpha > 1$, ρ takes its values in the interval $[1-1/\alpha,1/\alpha]$, the case $\rho = 1-1/\alpha$ corresponding to the spectrally positive situation ($\chi = 1$ and Z has no negative jumps) and the case $\rho = 1/\alpha$ to the spectrally negative situation ($\chi = -1$ and Z has no positive jumps). Of course these values coincide when $\alpha = 2$ (the Brownian case).

Because $\text{Re}(1/\Psi(\lambda) = O(|\lambda|^{-\alpha})$ as $|\lambda| \to \infty$, Theorem V.1 in Bertoin (1996) entails that Z possesses a local time process $L = \{L(t, x), t \ge 0, x \in \mathbb{R}\}$ in the sense that that for any non-negative Borel function f,

$$\int_0^t f(X_s) \, \mathrm{d}s = \int_{\mathbb{D}} f(x) L(t, x) \, \mathrm{d}x, \qquad t \ge 0.$$
 (2.2)

Besides, we know from Exercise V.3 in Bertoin (1996) that a.s. the map $x \mapsto L(t, x)$ is η -Hölder for every $\eta < (\alpha - 1)/2$. This property and the formula (2.2) allows to define the signed additive functionals

$$A_t^{(\beta)} = \int_0^t |Z_s|^\beta \operatorname{sgn}(Z_s) \, \mathrm{d}s, \quad t \ge 0,$$

for every $\beta > -(\alpha+1)/2$. For $\beta > -1$, the integral converges absolutely by the self-similarity of Z, because $\alpha > 1$. For $\beta \in (-(\alpha+1)/2, 1]$, it is meant as a Cauchy principal value:

$$A_t^{(\beta)} = \lim_{\varepsilon \to 0} \int_0^t \mathbf{1}_{\{|Z_s| > \varepsilon\}} |Z_s|^\beta \operatorname{sgn}(Z_s) \, ds = \lim_{\varepsilon \to 0} \int_{\mathbb{R}} \mathbf{1}_{\{|x| > \varepsilon\}} |x|^\beta \operatorname{sgn}(x) (L(t, x) - L(0, x)) \, dx.$$

We refer to the seminal paper Biane and Yor (1987) for numerous results on the principal values of Brownian local times, and also to Fitzsimmons and Getoor (1992) for their relations with certain limit theorems in the general stable case.

2.2. A symmetry lemma. Set $L_t = L(t, 0)$ for simplicity and let $\{\tau_t, t \geq 0\}$ be the right continuous inverse of L_t :

$$\tau_t = \inf \{ u \ge 0 : L_u > t \}, \quad t \ge 0.$$

From the strong Markov property for Z at times τ_t , and the fact that

$$\tau_{t+s} - \tau_t = \inf \{ u \ge 0 : L_{\tau_t+u} - t > s \}, \quad s, t \ge 0,$$

we see that the process $t \mapsto (\tau_t, A_{\tau_t}^{(\beta)})$ is a bivariate Lévy process. Besides, it is an easy consequence of the self-similarity of Z that $t \mapsto \tau_t$ is a stable subordinator with index $(\alpha-1)/\alpha$ and $t \mapsto A_{\tau_t}^{(\beta)}$ a stable process with index $(\alpha-1)/(\alpha+\beta)$. For the sake of concision, from now on we will use the notations $\gamma = (\alpha-1)/\alpha, \delta = (\alpha-1)/(\alpha+\beta)$, and set

$$\xi_t = A_{\tau_t}^{(\beta)}, \quad t \ge 0.$$

The following lemma, whose statement is trivial when Z itself is symmetric, entails that ξ is a symmetric δ -stable process:

Lemma 2.1. The random variables (τ_1, ξ_1) and $(\tau_1, -\xi_1)$ have the same law.

Proof. Introducing the process

$$Z_t^{\#} = \tau_1^{-1/\alpha} Z_{t\tau_1}, \qquad t \in [0, 1],$$

we see after a change of variable that

$$\xi_1 = \tau_1^{1+\beta/\alpha} \int_0^1 |Z_s^{\#}|^{\beta} \operatorname{sgn}(Z_s^{\#}) \, \mathrm{d}s.$$

On the other hand, it was established in Devulder et al. (2005) that for any bounded measurable functional F,

$$\mathbb{E}\left[F((Z_t^{\#}, \ t \in [0, 1]), \tau_1)\right] = \mathbb{E}^{(\text{br})}\left[\frac{c}{L_1}F\left((Z_t, \ t \in [0, 1]), L_1^{\alpha/(1-\alpha)}\right)\right], \quad (2.3)$$

where $\mathbb{P}^{(br)}$ stands for the law of the standard bridge associated with Z (see section VIII.3 in Bertoin (1996) for a definition) and $c \in (0, \infty)$ is a normalization constant. In passing, notice that (2.3) solves actually Exercise VIII.7 in Bertoin (1996) and corrects a typographical error therein: the density of the law of the

pseudo-bridge with respect to that of the standard bridge should be $kL(0,1)^{-1}$ and not $kL(0,1)^{1/(\alpha-1)}$.

Setting $\theta_1 = \tau^{-1-\beta/\alpha} \xi_1$, we get from (2.3) that for every f measurable bounded function

$$\mathbb{E}[f(\theta_1, \tau_1)] = \mathbb{E}^{(\mathrm{br})} \left[\frac{c}{L_1} f(A_1^{(\beta)}, L_1^{\alpha/(1-\alpha)}) \right].$$

But from their very definition,

$$(A_1^{(\beta)}, L_1^{\alpha/(1-\alpha)}) = (\bar{A}_1^{(\beta)}, \bar{L}_1^{\alpha/(1-\alpha)}),$$

where $\bar{A}^{(\beta)}$ and \bar{L} are associated with the time-reversed process $\bar{Z} = \{Z_{(1-t)-}, 0 \le t \le 1\}$. From Exercise VIII.5 in Bertoin (1996), this entails that

$$\mathbb{E}[f(\theta_1, \tau_1)] \ = \ \hat{\mathbb{E}}^{(\mathrm{br})} \left[\frac{c}{L_1} f(A_1^{(\beta)}, L_1^{\alpha/(1-\alpha)}) \right],$$

where $\hat{\mathbb{P}}^{(br)}$ is the law of the standard bridge associated with the dual process $\hat{Z} = -Z$. Using now (2.3) backwards, this yields

$$\mathbb{E}[f(\theta_1, \tau_1)] = \mathbb{E}[f(-\theta_1, \tau_1)],$$

for every f bounded measurable, so that (θ_1, τ_1) and $(-\theta_1, \tau_1)$ have the same law, as well as (ξ_1, τ_1) and $(-\xi_1, \tau_1)$.

Remarks 2.2. (a) This lemma is intuitively true from the facts that the variable (τ_1, ξ_1) is invariant under the operation reversing the time of each excursion of Z and making it right-continuous - see Definition (4.4) in Getoor and Sharpe (1981), and that the transformation of the law of Z under this latter operation yields actually the law of \hat{Z} - see Theorem (4.8) in Getoor and Sharpe (1981). Eventhough this argument is simpler, we feel that making it perfectly rigorous would force us to introduce some heavier notation than the one we use in the previous proof, and we leave this task to the interested reader.

(b) It is possible to determine the scaling parameter $\kappa_{\tau} = -\log \mathbb{E}[e^{-\tau_1}]$ of the subordinator τ from Proposition V.4 in Bertoin (1996). More precisely, we have

$$\kappa_{\tau} = 2\pi \left(\int_{\mathbb{R}} \operatorname{Re} \left(\frac{1}{1 + \Psi(\lambda)} \right) d\lambda \right)^{-1}$$

$$= \pi \kappa^{1/\alpha} \left(\int_{0}^{\infty} \left(\frac{1 + \lambda^{\alpha}}{(1 + \lambda^{\alpha})^{2} + \chi^{2} \tan^{2}(\pi\alpha/2)\lambda^{2\alpha}} \right) d\lambda \right)^{-1},$$

where κ is the scaling parameter of Z appearing in (2.1). When $\chi = 0$, using the 2nd formula p. 5 and the 3rd formula p. 10 in Magnus et al. (1966), one obtains

$$\kappa_{\tau} = \alpha \kappa^{1/\alpha} \sin(\pi/\alpha).$$

When $\chi \neq 0$, using the 5th formula p. 188 and the 8th formula p. 172 in Magnus et al. (1966), one gets

$$\kappa_{\tau} = \frac{\alpha \kappa^{1/\alpha} \sin(\pi/\alpha) \mu^{-1/\alpha} \sqrt{1 - \mu^2}}{\sin(\pi(\alpha - 1)|\rho - 1/2|) + \mu \sin(\pi|\rho - 1/2|)},$$

where $\mu = (1 + \chi^2 \tan^2(\pi \alpha/2))^{-1/2}$ and ρ is the positivity parameter of Z. Notice that when $|\chi| = 1$, the formula becomes

$$\kappa_{\tau} = \frac{\alpha \kappa^{1/\alpha}}{(\sin((\alpha - 1)\pi/2))^{1/\alpha}},$$

a fact which could have been obtained directly in solving Exercise VII.2 in Bertoin (1996).

However, it seems difficult to compute in general the scaling parameter $\kappa_{\xi} = -\log \mathbb{E}\left[e^{i\xi_1}\right]$ of the symmetric stable process ξ . In the Brownian case $\alpha = 2$, one can use the independence of the positive and negative excursions and solve an homogeneous Sturm-Liouville equation, which gives (see Formula (1 b) in Biane and Yor (1987))

$$\kappa_{\xi} = \frac{\pi (2\kappa)^{1/\delta - 1} 2^{\delta}}{2\delta \sin(\pi \delta/2)} \left(\frac{\delta^{\delta}}{\Gamma(\delta)}\right)^{2}.$$
 (2.4)

When $\alpha < 2$ one cannot use such a method because the excursions change a.s. their sign at least once during their lifetime. However, in the case $\beta = -1$, a remarkable formula due to Fitzsimmons, Getoor and Bertoin - see Theorem V.7 in Bertoin (1996) - entails that

$$\mathbb{E}\left[e^{\mathrm{i}\lambda A_{\tau_t}^{(-1)}}\right] = e^{-t\pi|\lambda|}, \quad \lambda \in \mathbb{R}$$

(notice that it is consistent with (2.4) above when $\delta=1$). The fact that κ_{ξ} does not depend on κ follows easily from the self-similarity of Z, and the simple constant π stems roughly from the identity $\mathcal{FH}=\pi i$ sgn where \mathcal{F} (resp. \mathcal{H}) is the Fourier (resp. Hilbert) transform. For $\beta=1$, using the identity $\mathcal{F}(x)=-i\,\mathcal{F}'$ and reasoning exactly as in section V.2 in Bertoin (1996), one can prove that the bivariate scaling parameter $\kappa(q,\lambda)=-\log\mathbb{E}[e^{-q\tau_1+i\lambda\xi_1}]$ appears in the linear equation

$$\lambda \varphi'(t) + \kappa(q,\lambda) = (q + \Psi(-t))\varphi(t)$$

satisfied by the function $\varphi = \mathcal{F}h$, where $T_{-x} = \inf\{t > 0, Z_t = -x\}$ and

$$h(x) = \mathbb{E}[e^{-qT_{-x} + i\lambda A_{T_{-x}}^{(1)}}]. \tag{2.5}$$

Using Lemma VIII.13 in Bertoin (1996), one can show that $|x|h(x) \in L^1(\mathbb{R})$, so that by (2.1) and the Riemann-Lebesgue lemma,

$$\kappa(q,\lambda) = \kappa(1 + \chi \tan(\pi \alpha/2)) \lim_{t \to +\infty} t^{\alpha} \varphi(-t).$$

However this formula seems barely tractable, even in the Brownian case where h is explicit - see Formula (5') p. 387 in Lachal (1991), but given as a quotient of Airy functions.

(c) When $\beta > -1$, the process ξ has finite variations and hence, can be decomposed as the difference of two i.i.d. δ -stable subordinators. There is actually another decomposition

$$\xi = \xi^+ - \xi^-,$$

with the notation

$$\xi_t^+ = \int_0^{ au_t} |Z_s|^{eta} \mathbf{1}_{\{Z_s > 0\}} \, \, \mathrm{d}s \quad ext{and} \quad \xi_t^- = \int_0^{ au_t} |Z_s|^{eta} \mathbf{1}_{\{Z_s < 0\}} \, \, \mathrm{d}s,$$

for every $t \geq 0$. Reasoning as above, one can show that the two δ -stable subordinators ξ^+ and ξ^- have also the same law. Nevertheless they are not independent since they jump simultaneously, unless $\alpha = 2$.

3. Proofs of the theorems

3.1. Proof of Theorem A. Let $\kappa(q,\lambda) = -\log \mathbb{E}[\mathrm{e}^{-q\tau_1 + \mathrm{i}\lambda\xi_1}]$ be the Laplace-Fourier exponent of the bivariate Lévy process $\{(\tau_t,\xi_t),\,t\geq 0\}$. Setting $I_t=\inf\{\xi_s,s\leq t\}$ and $S_t=\sup\{\xi_s,s\leq t\}$ for $t\geq 0$, the two-dimensional Wiener-Hopf factors are defined followingly:

$$\phi_+(q,\lambda) = \kappa(q,0) \int_0^\infty \mathbb{E}\left[\mathrm{e}^{-q au_t + \mathrm{i}\lambda S_t}
ight] \,\mathrm{d}t \quad ext{and} \quad \phi_-(q,\lambda) = \kappa(q,0) \int_0^\infty \mathbb{E}\left[\mathrm{e}^{-q au_t + \mathrm{i}\lambda I_t}
ight] \,\mathrm{d}t,$$

for all $q>0, \lambda\in\mathbb{R}$. The Wiener-Hopf factorization - see Theorem 1 in Isozaki (1996) - yields

$$\phi_{+}(q,\lambda)\phi_{-}(q,\lambda) = \frac{\kappa(q,0)}{\kappa(q,\lambda)}, \qquad q > 0, \lambda \in \mathbb{R}.$$

Notice that by Lemma 2.1, the Lévy processes (τ_t, ξ_t) and $(\tau_t, -\xi_t)$ have the same laws. This entails $\phi_-(q, \lambda) = \bar{\phi}_+(q, \lambda)$ for all q, λ , so that

$$|\phi_{+}(q,\lambda)| = \sqrt{\frac{\kappa(q,0)}{\kappa(q,\lambda)}}, \qquad q > 0, \lambda \in \mathbb{R}.$$
 (3.1)

The factor ϕ_+ is important for our problem. Indeed, setting

$$\Theta(r) = \inf\{t > 0, \, \xi_t > r\},\,$$

it follows from the Markov property at time $\Theta(r)$ - see Lemma 1 and Lemma 2 in Isozaki (1996) - that for every r>0

$$1 - \mathbb{E}\left[e^{-\tau_{\Theta(r)}}\right] = \int_{\mathbb{R}} \left(\frac{1 - e^{i\eta r}}{2\pi i\eta}\right) \phi_{+}(1,\eta) d\eta.$$

Since by scaling $\tau_{\Theta(r)} \stackrel{d}{=} r^{\frac{\alpha}{\alpha+\beta}} \tau_{\Theta(1)}$, after a change of variable we obtain

$$1 - \mathbb{E}\left[e^{-r\tau_{\Theta(1)}}\right] = \int_{\mathbb{R}} \left(\frac{1 - e^{i\eta}}{2\pi i\eta}\right) \phi_{+}(1, \eta r^{-(1+\beta/\alpha)}) d\eta, \qquad r \ge 0.$$
 (3.2)

We now proceed to an asymptotic analysis of $\phi_{+}(1,x)$ when $x \to \infty$, mimicking Lemma 3 in Isozaki (1996). First, by (3.1), a scaling argument and the symmetry of ξ ,

$$|\phi_{+}(1,x)| = \sqrt{\frac{\kappa(1,0)}{\kappa(1,x)}} = \sqrt{\frac{\kappa_{\tau}|x|^{-\delta}}{\kappa(|x|^{-\alpha/(\alpha+\beta)},1)}} \sim \sqrt{\frac{\kappa_{\tau}}{\kappa_{\xi}}}|x|^{-\delta/2}, \quad x \to \infty.$$
 (3.3)

Second, by Theorem 1 in Isozaki (1996) - noticing that $(e^{i\xi_t-1})$ should be read $(e^{i\xi_t}-1)$ therein, a scaling argument and a change of variable,

$$\arg(\phi_{+}(1,x)) = \int_{0}^{\infty} \frac{\mathrm{d}t}{t} \mathbb{E}\left[e^{-\tau_{t}} \sin(x\xi_{t}) \mathbf{1}_{\{\xi_{t}>0\}}\right]$$

$$= \operatorname{sgn}(x) \int_{0}^{\infty} \frac{\mathrm{d}t}{t} \mathbb{E}\left[e^{-t^{1/\gamma}\tau_{1}} \sin(|x|t^{1/\delta}\xi_{1}) \mathbf{1}_{\{\xi_{1}>0\}}\right]$$

$$= \delta \operatorname{sgn}(x) \int_{0}^{\infty} \frac{\mathrm{d}u}{u} \mathbb{E}\left[e^{-\tau_{1}\left(\frac{u}{|x|}\right)^{\alpha/(\alpha+\beta)}} \sin(|u|\xi_{1}) \mathbf{1}_{\{\xi_{1}>0\}}\right],$$

$$\sim \delta \operatorname{sgn}(x) \int_{0}^{\infty} \frac{\mathrm{d}u}{u} \mathbb{E}\left[\sin(|u|\xi_{1}) \mathbf{1}_{\{\xi_{1}>0\}}\right] = \frac{\delta \pi}{4} \operatorname{sgn}(x), x \to \infty(3.4)$$

Last, we compute an oscillating integral which will appear at the limit. Notice that this result is also quoted in Isozaki (1996) p. 223, nevertheless we will include a proof for the sake of clarity.

Lemma 3.1. One has

$$\int_{\mathbb{R}} \left(\frac{1 - \mathrm{e}^{\mathrm{i} \eta}}{2 \pi \mathrm{i} \eta} \right) |\eta|^{-\delta/2} \mathrm{e}^{\mathrm{i} \delta \mathrm{sgn}(\eta) \pi/4} \; \mathrm{d} \eta \; = \; \frac{2}{\delta \Gamma(\delta/2)} \, \cdot$$

Proof. Setting I for the value of the integral we can rewrite, by a parity argument,

$$I = \int_{\mathbb{R}} \left(\frac{\sin(\eta/2)}{\pi \eta} \right) |\eta|^{-\delta/2} e^{i(\delta \operatorname{sgn}(\eta)\pi/4 - \eta/2)} d\eta$$
$$= \frac{2}{\pi} \int_{0}^{\infty} \sin(\eta/2) \eta^{-1 - \delta/2} \cos(\delta \pi/4 - \eta/2) d\eta.$$

After trigonometric transformations and integrating by parts, we get

$$\begin{split} I &= \frac{\cos(\delta\pi/4)}{\pi} \int_0^\infty \eta^{-1-\delta/2} \sin\eta \, \mathrm{d}\eta \; + \; \frac{2\sin(\delta\pi/4)}{\pi\delta} \int_0^\infty \eta^{-\delta/2} \sin\eta \, \mathrm{d}\eta \\ &= \; \sin(\delta\pi/2) \left(\frac{\Gamma(1-\delta/2)}{\pi\delta} - \frac{\Gamma(-\delta/2)}{2\pi} \right) \; = \; \frac{2}{\delta\Gamma(\delta/2)}, \end{split}$$

where we used the 5th formula p. 9 in Magnus et al. (1966) in the second equality, and the formulae $\Gamma(z)\Gamma(-z)=-\pi/z\sin(\pi z), \Gamma(z)\Gamma(1-z)=\pi/\sin(\pi z)$ in the third.

Notice by the assumption on β , one has $r^{-(1+\beta/\alpha)} \to +\infty$ as $r \to 0$. Hence, putting (3.2), (3.3), (3.4) and Lemma 3.1 together one obtains

$$1 - \mathbb{E}\left[e^{-r\tau_{\Theta(1)}}\right] \sim \sqrt{\frac{\kappa_{\tau}}{\kappa_{\xi}}} \left(\frac{2}{\delta\Gamma(\delta/2)}\right) r^{\gamma/2}, \qquad r \to 0,$$

the dominated convergence argument being plainly justified by (3.2) and (3.3). By a Tauberian theorem and a monotone density theorem - see e.g. Bertoin (1996) p. 10, this yields

$$\mathbb{P}[\tau_{\Theta(1)} > t] \sim \sqrt{\frac{\kappa_{\tau}}{\kappa_{\xi}}} \left(\frac{2}{\delta \Gamma(\delta/2) \Gamma(1 - \gamma/2)} \right) t^{-\gamma/2}, \qquad t \to +\infty.$$
 (3.5)

The key-point is now that by the definition of T, for every t > 0

$$\{T > t\} = \{A_s^{(\beta)} < 1, \ s \le t\} \subset \{\xi_u < 1, \ \forall \ u \ / \ \tau_u \le t\} = \{\Theta_1 > L_t\} \subset \{\tau_{\Theta(1)} > t\}.$$

This entails finally

$$\mathbb{P}[T > t] \le \mathcal{K} t^{-\gamma/2}, \quad t \to +\infty,$$

with

$$\mathcal{K} \; = \; \sqrt{\frac{\kappa_{\tau}}{\kappa_{\xi}}} \left(\frac{2}{\delta \Gamma(\delta/2) \Gamma(1 - \gamma/2)} \right)$$

explicit when $\alpha = 2$ or $\beta = -1$, thanks to the above Remark 2 (a).

Remark 3.2. At a less precise level and for the case where $\beta = 1$ and Z is symmetric, it had been shown in Devulder et al. (2005) that

$$\mathbb{P}\left[\sup_{t\in[0,1]}A_t^{(1)}<\varepsilon\right] \leq \varepsilon^{(\alpha-1)/2(\alpha+1)+o(1)}, \qquad \varepsilon\to 0,$$

with a different method: first, one remarks that

$$\mathbb{P}\left[\sup_{t\in[0,\tau_1]}A_t^{(1)}<\varepsilon\right]\leq \mathbb{P}\left[\sup_{t\in[0,1]}A_{\tau_t}^{(1)}<\varepsilon\right]\sim k\varepsilon^{(\alpha-1)/2(\alpha+1)}, \qquad \varepsilon\to 0,$$

for some finite positive constant k, because $t \mapsto A_{\tau_t}^{(1)}$ is a symmetric $(\alpha - 1)/(\alpha + 1)$ -stable process and in view of Proposition VIII.3 in Bertoin (1996). Second, at the cost of some o(1) term on the exponent, one replaces τ_1 by 1 in the probability on the left-hand side after a suitable slicing of the big values of τ_1 .

3.2. Proof of Theorem B. Set $\Theta = \inf\{t > 0, \, \xi_t \geq 1\}$ and for every c > 0 introduce the function

$$f^{c}(t) = (c \log t)^{1/(1-\alpha)}, \qquad t > 1.$$

One has

$$\begin{split} \mathbb{P}\left[\Theta > t^{\gamma}\right] & \leq & \mathbb{P}\left[\tau_{t^{\gamma}/2} < tf^{c}(t)\right] + \mathbb{P}\left[\Theta > t^{\gamma}, \, \tau_{t^{\gamma}/2} \geq tf^{c}(t)\right] \\ & \leq & \mathbb{P}\left[\tau_{t^{\gamma}/2} < tf^{c}(t)\right] + \mathbb{P}\left[\tau_{\Theta^{-}} \geq tf^{c}(t)\right], \end{split}$$

since τ is a.s. increasing. The key-point is the inequality

$$T \geq \tau_{\Theta-}$$
 a.s. (3.6)

Indeed, if we set $T'=\inf\{t\geq T,\, Z_t=0\}$, then it follows from the absence of negative jumps for Z that $Z_T\geq 0$ and $A_{T'}^{(\beta)}\geq 1$ a.s. Hence, since a.s. $T\leq \tau_\Theta$ we see that a.s. $T'=\tau_\Theta$, which readily entails (3.6) by the definition of T'. We deduce

$$\begin{split} \mathbb{P}\left[T \geq t f^c(t)\right] & \geq & \mathbb{P}\left[\Theta > t^{\gamma}\right] - \mathbb{P}\left[\tau_{t^{\gamma}/2} < t f^c(t)\right] \\ & \geq & \mathbb{P}\left[\Theta > t^{\gamma}\right] - \mathbb{P}\left[\tau_{1/2} < f^c(t)\right] \\ & \geq & k_1 t^{-\gamma/2} - t^{-ck_2}, \qquad t \to \infty, \end{split}$$

for some $k_1, k_2 > 0$, where in the third inequality we used Proposition VIII. 2 and the remark after Proposition VIII.3 p. 221 in Bertoin (1996) - recall that τ is a γ -stable subordinator. Taking c sufficiently large, we obtain

$$\mathbb{P}\left[T > t f^{c}(t)\right] > k_{3} t^{-\gamma/2}, \qquad t \to \infty.$$

for some positive constant k_3 . Changing the variable $u = tf^c(t)$, we see by the definition of f^c that $t \leq u(\log u)^{1/(\alpha-1)}/2$ provided u is sufficiently large. Putting everything together yields

$$\mathbb{P}\left[T \ge u\right] \ge k_4 u^{-\gamma/2} (\log u)^{-1/2\alpha}, \qquad u \to \infty,$$

for some constant $k_4 > 0$, as required.

Remarks 3.3. (a) In the above proof, the key-argument (3.6) is completely peculiar to the spectrally positive framework. When there are negative jumps, it seems possible that

$$T' \ll \tau_{\Theta}$$

with significant probability, and this raises the question whether the exponent $\gamma/2$ is still the critical one. In Simon (2007) we actually proved that it is not the case anymore when there are negative jumps, at least when α is close to 1, hence contradicting the validity of Shi's conjecture in general. At the end of this paper, we

will state another conjecture on the value of the critical exponent θ for integrated Lévy α -stable processes.

(b) In the Brownian case $\alpha = 2$, it is possible to obtain a sharper bound

$$\mathbb{P}[T > t] > \mathcal{K}' t^{-1/4}, \qquad t \to \infty,$$

for some explicit constant \mathcal{K}' - see Theorem 3 in Isozaki (1996). The method consists in proving that

$$\mathbb{P}\left[\tau_{\Theta} - T > t\right] \leq \mathcal{K}'' t^{-1/4}, \qquad t \to \infty,$$

where \mathcal{K}'' is made explicit with the help of the function h in (2.5), and shown to be strictly smaller than \mathcal{K} . Notice that the arguments of Isozaki (1996) are given only for $\beta > 0$. Nevertheless they can be extended to the general situation $\beta > -3/2$ with the same formula for \mathcal{K} and \mathcal{K}' . As mentioned in the introduction, in a subsequent work Isozaki and Kotani (2000) it is proved that

$$\mathbb{P}[T > t] \sim \mathcal{K}_{\beta} t^{-1/4}, \qquad t \to \infty$$

as soon as $\beta > -1$, for some explicit constant \mathcal{K}_{β} - see Corollary 1 in Isozaki and Kotani (2000). Notice that the value $\mathcal{K}_1 = 3^{4/3}\Gamma(2/3)/\pi 2^{13/12}\Gamma(3/4)$ had been previously computed by Goldman (1971).

(c) In the Non-Gaussian case, it seems difficult to obtain sharper estimates with the above method, due to the lack of information on h. When $\beta \geq 0$ however, it is possible to obtain an explicit uniform bound

$$\mathbb{P}\left[\tau_{\Theta} - T > t\right] \leq \mathcal{K} t^{-\gamma/2}, \qquad t \geq 0, \tag{3.7}$$

thanks to the formula

$$\mathbb{P}\left[\tau_{\Theta} - T > t\right] = \alpha \int_0^\infty \widehat{p}_1(u) \, \mathbb{P}[Z_T > t^{1/\alpha} u] \, \mathrm{d}u, \tag{3.8}$$

where \hat{p}_1 is the density of $\hat{Z}_1 = -Z_1$, and an upper bound on $\mathbb{P}[Z_T > u]$ which will be detailed soon afterwards. The formula (3.8) follows at once from the Markov property at time T, an integration by parts, a scaling argument and the well-known fact that

$$\sup\{\widehat{Z}_t, t \le 1\} \stackrel{d}{=} \widehat{Z}_1 \mathbf{1}_{\{\widehat{Z}_1 > 0\}}, \tag{3.9}$$

since \widehat{Z} has no positive jumps - see e.g. Bingham (1975) p. 749. Unfortunately, it seems that (3.7) is useless as soon as we do not have any more precise information on κ_{ξ} .

- (d) Recently, the simple method of Theorem B has been applied successfully in Lachal and Simon (2007) to derive some a priori difficult lower bound on the small ball probabilities for a class of homogeneous Brownian functionals.
- 3.3. Final remarks and conjectures. In this paragraph we first give some comments on the positive moments of the interesting random variable Z_T . In the Brownian case and for $\beta=1$, the law of Z_T had been computed by Gor'kov see Formula (3) in Lachal (1991), but the expression seems too complicated to allow further computations. In the general case with no negative jumps, using (3.8), the formula

$$\mathbb{E}[S^k] = k \int_0^\infty s^{k-1} \mathbb{P}[S > s] \, \mathrm{d}s,$$

which is valid for any positive random variable S and any k>0, and Fubini's theorem, one obtains

$$\mathbb{E}[(\tau_{\Theta} - T)^{k}] = \alpha \left(\int_{0}^{\infty} u^{-k\alpha} \widehat{p}_{1}(u) du \right) \mathbb{E}[Z_{T}^{k\alpha}]$$
$$= \left(\frac{\Gamma(1 - k\alpha)}{\kappa^{k} \Gamma(1 - k)} \right) \mathbb{E}[Z_{T}^{k\alpha}],$$

for every $k \in [0, 1/\alpha)$, where the computation of the integral in the second line comes from Theorem 2.6.3 in Zolotarev (1986). In the Brownian case, using (5.6) in Isozaki (1996) yields then the following criterion:

$$k < 1/2 \iff \mathbb{E}[Z_T^k] < +\infty,$$

for every k > 0. In the general stable case with no negative jumps, using Theorem A we obtain the inclusion

$$k < (\alpha - 1)/2 \quad \Longrightarrow \quad \mathbb{E}[Z_T^k] < +\infty, \tag{3.10}$$

which we actually believe to be an equivalence. Notice that when $\beta \geq 0$, one can also obtain the following reinforcement of (3.10):

$$\mathbb{P}[Z_T > u] \leq \mathcal{K} u^{-(\alpha - 1)/2}, \quad u \to \infty,$$

for some explicit constant K, which together with Theorem 2.6.3 in Zolotarev (1986) leads to the aforementioned estimate (3.7). Let me sketch the proof. First, by the Markov property at time T, a scaling argument and Skorokhod's formula (3.9) one gets

$$\mathbb{P}[Z_T > u] = \frac{\mathbb{P}[Z_T > u, T_{cZ_T} \circ \theta_T > Z_T^{\alpha}]}{\alpha \mathbb{P}[0 < \widehat{Z}_1 < 1 - c]},$$

for every $c \in (0,1)$ and u > 0, where θ is the shift operator of Z and T_x the first hitting time of x by Z. Second, since Z has no negative jumps and since $\beta \geq 0$, one deduces

$$\mathbb{P}[Z_T > u] \le \frac{\mathbb{P}[A_{\Theta(1)}^{(\beta)} > cu^{\alpha+\beta}]}{\alpha \mathbb{P}[0 < \widehat{Z}_1 < 1 - c]}$$

for every $c \in (0,1)$ and u > 0. Last, remembering that $A^{(\beta)}$ is a symmetric δ -stable process and using Exercise VIII. 3 in Bertoin (1996), one obtains

$$\mathbb{P}[Z_T > u] \leq \mathcal{K}_c u^{-(\alpha - 1)/2}, \quad u \to \infty,$$

for an explicit constant \mathcal{K}_c which can be minimized in $c \in (0,1)$.

In the general case with or without negative jumps, finding a lower bound on the quantities $\mathbb{P}[Z_T > u]$ seems at least as difficult as obtaining a lower bound on $\mathbb{P}[T > u]$. In the case without *positive* jumps, the two asymptotics are clearly connected thanks to the basic inequality

$$\mathbb{P}[Z_T > u] \le \mathbb{P}[T > u^{\alpha(1-o(1))}] + \mathbb{P}[S_1 > u^{o(1)}], \quad u \to \infty,$$

and the fact that $S_1 = \sup\{Z_t, t \leq 1\}$ has finite exponential moments - see Corollary VII.2 in Bertoin (1996). From the Lévy-Khintchine formula and the optional sampling theorem at time T, one can finally raise the following

Conjecture. For every k > 0 one has

$$\mathbb{E}[T^k] < +\infty \quad \Longleftrightarrow \quad \mathbb{E}[Z_T^{k\alpha}] < +\infty.$$

To put it in a nutshell, in this paper we have showed that the critical value of k in the above conjecture is $(\alpha - 1)/2\alpha$ when there are no negative jumps, but we still do not know what this value should be when there are negative jumps. It might be worth mentioning that if we set $S = \inf\{t > 0, Z_t > 1\}$, then from Proposition VIII. 2 and Exercise VIII. 3 in Bertoin (1996) we have

$$k < \rho \iff \mathbb{E}[S^k] < +\infty \iff \mathbb{E}[Z_S^{k\alpha}] < +\infty,$$

as soon as Z has positive jumps (when Z has no positive jumps only the first equivalence remains valid since $Z_S = 1$ a.s.) From this fact and the results of the present article, one might suppose that in general the critical exponent is

$$k_0 = \frac{\rho}{2} \tag{3.11}$$

in the above conjecture - recall that $\rho=(\alpha-1)/\alpha$ when Z has no negative jumps and $\alpha>1$. Besides, when $\beta=1$, one can show from Proposition 3.4.1 in Samorodnitsky and Taqqu (1994) that

$$\mathbb{P}[A_1^{(1)} > 0] = \rho,$$

and it would be somewhat surprising that this quantity does not play any rôle in the lower tails of $A^{(1)}$, since it does in the lower tails of Z.

Let us finally mention that if (3.11) held, then it would also be very interesting to know if for $\beta = 1$ the following upper bound dealing with the *non-strictly* stable process $Z' = \{Z_t - t, t \ge 0\}$, remains true:

$$\mathbb{P}\left[\int_0^t Z_s \, \mathrm{d}s < \varepsilon + t^2, \, \forall \, t \in [0, 1]\right] \leq \varepsilon^{\rho/2 + o(1)}, \qquad \varepsilon \to 0.$$
 (3.12)

Indeed, following a seminal approach due to Sinai (1992b), Molchan and Kholkhov recently showed that upper estimates like (3.12) provide the key-step to get upper bounds on the Hausdorff dimension of the so-called regular Lagrangian points for the inviscid Burgers equation with random initial data - see Theorem 1 in Molchan and Khokhlov (2004). Molchan and Kholkhov's arguments are given in a fractional Brownian framework, but one can show Simon (2007) that they remain valid for Lévy stable processes: if (3.12) held, then one would have

$$\operatorname{Dim}_H \mathcal{L} < 1 - \rho < 1/\alpha$$
 a.s.

where \mathcal{L} is the set of Lagrangian regular points for the inviscid Burgers equation with initial data \widehat{Z} . We refer to Sinai (1992b); Bertoin (1998); Janicki and Woyczynski (1997) for various presentations of the Burgers equation with random initial impulse and we recall that in Janicki and Woyczynski (1997), the following well-known conjecture is made:

$$Dim_H \mathcal{L} = 1/\alpha \quad a.s. \tag{3.13}$$

This conjecture is true for Brownian motion by the result of Sinai (1992b) - see also Bertoin (1998) for more rigorous arguments and more general results, and when \widehat{Z} has no positive jumps i.e. Z has no negative jumps and $\rho=1-1/\alpha$ - see again Bertoin (1998). The situation where \widehat{Z} has positive jumps seems to require a completely different methodology - see the conclusion of Bertoin (1998), which we believe to be connected with the optimal upper bound on $\mathbb{P}[T>u]$ when Z has

negative jumps: in Simon (2007), we proved that in this situation (3.13) is false at least when α is close to 1, and we conjectured that

$$\operatorname{Dim}_H \mathcal{L} = 1 - \rho$$
 a.s.

with the above notations, which is clearly connected to (3.11) as far as the upper bound is concerned.

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